# EE 508

### Lecture 36

# Oscillators, VCOs, and Oscillator/VCO-Derived Filters

### Review from last lecture:

Only two of these circuits are useful directly as bias generators!



#### Transconductance Linearization Strategies Review from last lecture:



### Programmable Filter Structures Review from last lecture:

It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage





Filter characteristics can be controlled by an analog voltage ( $V_{CTRL}$ ) an analog current ( $I_{CTRI}$ ), or a Boolean signal

Much more controllability than with a potentiometer

It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage





It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage



$$
V_{OUT} = \frac{s^2 C_1 C_2 V_C + s C_1 g_{m2} V_B + g_{m1} g_{m2} V_A}{s^2 C_1 C_2 + s C_1 g_{m2} + g_{m1} g_{m2}} \qquad \omega_0 = \sqrt{\frac{g_{m1} g_{m2}}{C_1 C_2}}
$$



Programmable Second-Order Filter

It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage



Programmable Second-Order Filter

It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage



Programmable Filter Structures

\nbe assumed that the transconductance gain can be programmed with

\n
$$
V_{A} = \frac{g_{\text{max}}}{c_{1}} \frac{g_{\text{max}}}{c_{2}} \frac{g_{\text{max}}}{c_{2}} \frac{g_{\text{max}}}{c_{2}} \frac{g_{\text{max}}}{c_{2}}
$$
\n
$$
V_{\text{out}}
$$
\n
$$
V_{\text{OUT}} = \frac{s^{2}C_{1}C_{2}V_{C} + sC_{1}g_{\text{max}}V_{B} + g_{\text{max}}g_{\text{max}}V_{A}}{s^{2}C_{1}C_{2} + sC_{1}g_{\text{max}}V_{B} + g_{\text{max}}g_{\text{max}}}
$$
\nProgrammable Second-Order Filter

\nProofa

\n
$$
V_{\text{out}}
$$
\n
$$
V_{\text{out
$$

It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage



 $0 \leq l \leq \alpha$ 

 $1\overline{\smash{0}}\,2$ 

 $=\sqrt{\frac{c_2}{c}}\frac{\sqrt{5m15m2}}{c}$ 

 $g_{\rm ml} g_{\rm m2}$ 

 $C_1$  C<sub>2</sub>

$$
V_{OUT} = \frac{s^2 C_1 C_2 V_C + s C_1 g_{m3} V_B + g_{m1} g_{m2} V_A}{s^2 C_1 C_2 + s C_1 g_{m3} + g_{m1} g_{m2}} \qquad \omega_0 = \sqrt{\frac{g_{m1} g_{m2}}{C_1 C_2}}
$$
  
Programmable Second-Order Filter  

It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage





# Programmable Filter Components

It will be assumed that the transconductance gain can be programmed with either a dc current or a dc voltage



If  $\mathsf{Z}_\mathsf{L}$  is a capacitor serve as either positive or negative programmable inductors

**Many other useful programmable filter components and filter structures possible**

# Question:

What is the relationship, if any, between a filter and an oscillator or VCO?



i.e. Can an oscillator be viewed as a filter with no input?

### What is the relationship, if any, between a filter and an oscillator or VCO?



Will focus on modifying oscillator structures to form high frequency narrowband filters

Claim: Narrow band filters are dependent primarily on the poles close to the imaginary axis and affected little by poles that are farther away

Goal: Obtain very high frequency filter structures

### What is the relationship, if any, between a filter and an oscillator or VCO?





- When power is applied to an oscillator, it initially behaves as a small-signal linear network
- When operating linearly, the oscillator has poles (but no zeros)
- Poles are ideally on the imaginary axis or appear as cc pairs in the RHP
- There is a wealth of literature on the design of oscillators
- Oscillators often are designed to operate at very high frequencies
- If cc poles of a filter are moved to RHP it will become an oscillator
- If RHP poles of an oscillator are moved to the LHP it will become a filter if an input is applied **Scillator**<br>
• When power is applied to an oscillator, it initially behaves as a small-<br>
• When operating linearly, the oscillator has poles (but no zeros)<br>
• Poles are ideally on the imaginary axis or appear as cc pairs
- 

# Oscillator Background:

Consider a cascaded integrator loop comprised of n integrators

This structure is often used to build oscillators





(assume an odd number of inverting integrators)

$$
X_{OUT} = -\left(\frac{I_0}{s}\right)^n X_{OUT}
$$

$$
X_{OUT}(s^n + I_0^n) = 0
$$

$$
D(s) = s^n + I_0^n
$$

Consider the poles of  $\mathsf{D}(\mathsf{s})\,{=}\,\mathsf{s}^{\mathsf{n}}\texttt{+}I_0^{\mathsf{n}}$  $=$  S<sup>''</sup>+  $I$ 

1  $n + I_0^n = 0$  $\mathsf{S}^{\prime\prime}$ +  $I_{0}^{\prime\prime}$  = n n  $\textbf{S}^{\prime \prime} = -I_0^\prime$ n  $\left\vert \mathbf{S}\right\vert =\right\vert -I_{0}^{\mathsf{H}}\left\vert ^{n}\right\vert ^{n}$  $=\left[-I_0^n\right]$ 

 $\left[-1\right]$ 

 $\mathbf{S} = \begin{bmatrix} -1 \end{bmatrix}$ n  $\begin{bmatrix} I_0^{\mathsf{H}} \end{bmatrix}$ n  $=-\left[-1\right]_n^{\frac{1}{n}}\left[I_0^n\right]$ 

 $\textsf{S}=I_0$   $\mid$   $\mid$   $\mid$   $\mid$   $\mid$ 

 $\left[-1\right]$ 

 $\lfloor \ln \lfloor \ln \right]$ 

 $1 - 1$ 

1

Poles are the n roots of -1 scaled by  $I_0$ 

Roots of -1:



Roots are uniformly spaced on a unit circle

Consider the poles of  $\,\,\mathsf{D}\big(\,\mathsf{S}\big) \!=\! \,\mathsf{S}^{\mathsf{n}} \texttt{+} \, I_{0}^{\mathsf{n}}$  $=$  S<sup>''</sup>+  $I$ 















### Some useful theorems

Theorem: A rational fraction 
$$
T(s) = \frac{N(s)}{\prod_{i=1}^{n} (s - p_i)}
$$
  
\nin partial fraction form as  $T(s) = \sum_{i=1}^{n} \frac{A_i}{s - p_i}$   
\nwhere  $A_i = (s - p_i)T(s)|_{s = p_i}$  for  $1 \le j \le n$ 

 $\mathbf{h}(\mathrm{t})$ = $\sum \mathrm{A_{i}}\mathrm{e}^{\mathrm{p_{i}t}}$  where the  $\sum_{k=1}^{n} A_k \rho_k$  and  $\sum_{k=1}^{n} A_k$ i<sup>u</sup> n i=1 r(t)= $\sum A_i e^{p_i t}$  where the numbers  $A_i$  are the Theorem: The impulse response of a rational fraction T(s) with simple poles can be expressed as  $\;\;$  r $(t)$ = $\stackrel{..}{\Sigma}$  A ${_{\rm i}}$ e $^{\rm p,t}$  where the numbers A ${_{\rm i}}$  are the coefficients in the partial fraction expansion of  $T(s)$ 

Theorem: If  $p_i$  is a simple complex pole of the rational fraction T(s), then the partial fraction expansion terms in the impulse response corresponding to  $\boldsymbol{\mathsf{p}}_\mathsf{i}$  and  $\boldsymbol{\mathsf{p}}_\mathsf{i}^\star$ can be expressed as  $A_i^*$  ,  $A_i^*$ i\*s-p<sub>i</sub> s-p<sub>i</sub> A A  $^+$ 

Theorem: If  $p_i = \alpha_i + j\beta_i$  is a simple pole with non-zero imaginary part of the rational fraction T(s), then the impulse response terms corresponding to the poles  $\mathsf{p_i}$  and  $\mathsf{p_i}^{\star}$ in the partial fraction expansion can be expressed as

$$
A_i|e^{\alpha_i t} \cos(\beta_i t + \theta_i)
$$

where  $\, \theta_{\text{\tiny{i}}}$  is the angle of the complex quantity  $\mathsf{A}_{\text{\tiny{i}}}$ 

Observe r(t) term corresponding to any complex pole pair is real !

Theorem: If all poles of an n-th order rational fraction T(s) are simple and have a non-zero Imaginary part, then the impulse response can be expressed as

$$
\sum_{i=1}^{n/2} |A_i| e^{\alpha_i t} \cos(\beta_i t + \theta_i)
$$

where  $\ \theta_{\text{i}}$  ,A $_{\text{i},}$ α $_{\text{i}}$ , and β $_{\text{i}}$  are as defined before

i

Theorem: If an odd-order rational fraction has one pole on the negative real axis at  $\alpha_0$  and n simple poles that have a non-zero Imaginary part, then the impulse response can be expressed as

$$
A_0e^{\alpha_s t}+\sum_{i=1}^{n/2}|A_i|e^{\alpha_i t}cos(\beta_i t+\theta_i)
$$

where  $\ \theta_{\text{i}}$  ,A $_{\text{i},}$ α $_{\text{i}}$ , and β $_{\text{i}}$  are as defined before

Observe r(t) is real for both even and odd n !

### Poles of  $\mathsf{D}(\mathsf{s})\,{=}\,\mathsf{s}^{\mathsf{n}}\texttt{+}I_0^{\mathsf{n}}$  $=$  S<sup>''</sup>+  $I$ Consider the following 3-pole situation

![](_page_22_Figure_1.jpeg)

 $\varepsilon \bullet |A_i| e^{\alpha_i t} \cos(\beta_i t + \theta_i)$ Oscillatory output at startup with any small impulse input  $\, \varepsilon \, \partial(t)$  :

Starts at  $\omega = \beta = 0.866I_0$  and will slow down as nonlinearities limit amplitude

### Consider the following 3-pole situationPoles of  $\quad \mathsf{D}(\mathsf{s})\,{=}\,\mathsf{s}^{\mathsf{n}}\texttt{+}I_0^{\mathsf{n}}$  $=$  S<sup>''</sup>+  $I$ Consider moving all poles to left by Δα  $n=3$  $\mathsf{Im}$  $β=0.866 I_0$ X  $\alpha$ =0.5 I<sub>0</sub> - Δα

![](_page_23_Figure_1.jpeg)

So since α is fixed , to get a high  $ω_0$ , want β as large as possible

![](_page_24_Picture_0.jpeg)

# Stay Safe and Stay Healthy !

# End of Lecture 36

### **Consider now the filter obtained by adding a loss of α<sub>L</sub> to the integrators**

Will now determine  $\alpha_{\text{l}}$  and  $I_0$  needed to get a desired pole Q and  $\omega_0$  by moving all poles so that right-most pole pair is the dominant high-frequency pole pair of the filter

The values of  $\alpha$  and  $\beta$  are dependent upon  $I_0$  but the angle θ is only dependent upon the number of integrators in the oscillator or VCO

 $\alpha$ +j $\beta$   $=$   $I_{0} \big( cos\theta + j sin\theta \big)$ 

Define the location of the filter pole to be

$$
\alpha_F\text{+}j\beta_F
$$

It follows that

$$
\beta_F = \beta \qquad \alpha_F = \alpha - \alpha_L
$$

The relationship between the filter parameters is well known

$$
\beta_F = \frac{\omega_0}{2Q} \sqrt{4Q^2 - 1} \qquad \alpha_F = -\frac{\omega_0}{2Q}
$$

Thus for any n

$$
I_0 = \frac{\omega_0}{(\sin \theta) 2Q} \sqrt{4Q^2 - 1} \qquad \alpha_L = \frac{\omega_0}{2Q} + I_0 \cos \theta = \frac{\omega_0}{2Q} + \frac{\omega_0}{2Q(\tan \theta)} \sqrt{4Q^2 - 1}
$$

![](_page_26_Figure_12.jpeg)

**Will a two-stage structure give the highest frequency of operation for integrators with unity gain frequency I0?**

![](_page_27_Figure_1.jpeg)

$$
\omega_0 = \sqrt{(\alpha - \Delta \alpha)^2 + \beta^2} \quad \longrightarrow \quad \omega_0 = \sqrt{(-\Delta \alpha)^2 + \beta^2}
$$

- Even though the two-stage structure may not oscillate, can work as a filter!
- Need odd number of inversions in integrators
- Can add phase lead if necessary

# Oscillator Background:

What will happen with a circuit that has two pole-pairs in the RHP?

![](_page_28_Figure_2.jpeg)

General form of response for odd number of poles:

$$
A_0e^{\alpha_{\scriptscriptstyle 0} t}+\sum_{i=1}^{n/2}\big|A_i\big|e^{\alpha_{\scriptscriptstyle i} t}cos\big(\beta_{i}t+\theta_{i}\big)
$$

The impulse response (for n=7) will have two decaying exponential terms  $A_0e^{\alpha_0 t} + \sum_{i=1}^{\infty} |A_i|e^{\alpha_i t} \cos(\beta_i t + \theta_i)$ <br>The impulse response (for n=7) will have two decaying exponential terms

#### **What will happen with a circuit that has two pole-pairs in the RHP?**

![](_page_29_Figure_1.jpeg)

![](_page_29_Picture_158.jpeg)

$$
\alpha_1 = 0.2225
$$
  $\beta_1 = 0.974$ 

 $\alpha_2 = 0.9009$   $\beta_2 = 0.4338$ 

Consider the growing exponential terms and normalize to  $I_0=1$ 

$$
|A_1|e^{\alpha_1 t}\cos(\beta_1 t + \theta_1) + |A_2|e^{\alpha_2 t}\cos(\beta_2 t + \theta_2)
$$

At t=145 (after only 10 periods of the lower frequency signal)

$$
r = \frac{e^{\alpha_{2}t}}{e^{\alpha_{1}t}}\Big|_{t=145} = \frac{e^{.9009 \bullet 145}}{e^{.2225 \bullet 145}} = 5.2 \times 10^{42}
$$

### **The lower frequency oscillation will completely dominate !**

#### **What will happen with a circuit that has two pole-pairs in the RHP?**

![](_page_30_Figure_1.jpeg)

Thanks to Chen for these plots

Figure 14 N=8 impulse response

### **Can only see the lower frequency component !**

#### **What will happen with a circuit that has two pole-pairs in the RHP?**

![](_page_31_Figure_1.jpeg)

![](_page_31_Figure_2.jpeg)

# **After even only two periods of the lower frequency waveform, it**

**How do we guarantee that we have a net coefficient of +1 in D(s)?**

$$
\mathsf{D}(\mathsf{s})\,{=}\,\mathsf{s}^{\mathsf{n}}\text{+}\,I_0^{\mathsf{n}}
$$

![](_page_32_Figure_2.jpeg)

Must have an odd number of inversions in the loop !

If n is odd, all stages can be inverting and identical !

**How do we guarantee that we have a net coefficient of +1 in D(s)?**

$$
\mathsf D(\mathsf s)\,{=}\,\mathsf s^{\mathsf n}\text{-}\,I_0^{\mathsf n}
$$

![](_page_33_Figure_2.jpeg)

If fully differential or fully balanced, must have an odd number of crossings of outputs

Applicable for both even and odd order loops

# Inputs to Oscillator-Derived Filters:

Most applicable to designing 2<sup>nd</sup>-order high frequency narrow band bandpass filters

- Add loss to delay stages
- Multiple Input Locations Often Possible
- Natural Input is Input to delay stage

![](_page_34_Figure_5.jpeg)

![](_page_34_Figure_6.jpeg)

### **A lossy integrator stage**

![](_page_35_Figure_1.jpeg)

$$
I(s) = \frac{-g_{m1}/C_X}{s + g_{m2}/C_X}
$$

$$
I_0 = g_{m1}/C_X
$$

$$
\alpha_L = g_{m2}/C_{\chi}
$$

### **A fully-differential voltage-controlled integrator stage**

![](_page_36_Figure_1.jpeg)

Will need CMFB circuit

### **A fully-differential voltage-controlled integrator stage with loss**

![](_page_37_Figure_1.jpeg)

### **A fully-differential voltage-controlled integrator stage with loss**

(almost same as previous)

![](_page_38_Figure_2.jpeg)

Using the single-stage lossy integrator, design the integrator to meet a given  $\omega_0$  and Q requirement  $V_{DD}$ 

Recall:

$$
I_0 = \frac{\omega_0}{(\sin \theta) 2Q} \sqrt{4Q^2 - 1} \qquad (1)
$$
  

$$
\alpha_L = \frac{\omega_0}{2Q} + \frac{\omega_0}{2Q(\tan \theta)} \sqrt{4Q^2 - 1} \qquad (2)
$$

Substituting for  $I_0$  and  $\alpha$ <sub>L</sub> we obtain:

$$
\frac{g_{m1}}{C_{x}} = \frac{\omega_{0}}{(sin\theta)2Q} \sqrt{4Q^{2}-1}
$$
\n
$$
\frac{g_{m2}}{C_{x}} = \frac{\omega_{0}}{2Q} + \frac{\omega_{0}}{2Q(tan\theta)} \sqrt{4Q^{2}-1}
$$
\n(3)\n
$$
\omega_{L} = g_{m2}/C_{x}
$$
\n
$$
\omega_{L} = g_{m2}/C_{x}
$$

![](_page_39_Figure_6.jpeg)

$$
I_0 = g_{m1}/C_{\chi}
$$

$$
\alpha_L = g_{m2}/C_{\chi}
$$

Unknowns:  $I_B, V_{EB1}, W_1/L_1, W_2/L_2, C_X$ 

Using the single-stage lossy integrator, design the integrator to meet a given  $\omega_0$  and Q requirement  $V_{DD}$ 

(5)

Expressing  $g_{m1}$  and  $g_{m2}$  in terms of design parameters:<br> $\frac{1}{\sqrt{1-\frac{1}{$ 

$$
\frac{\mu C_{OX} W_1 V_{EB1}}{L_1 C_X} = \frac{\omega_0}{(\sin \theta) 2Q} \sqrt{4Q^2 - 1}
$$

$$
\frac{\mu C_{OX} W_2 V_{EB2}}{L_2 C_X} = \frac{\omega_0}{2Q} + \frac{\omega_0}{2Q(\tan\theta)} \sqrt{4Q^2 - 1}
$$
 (6)

![](_page_40_Figure_5.jpeg)

$$
I_0 = g_{m1}/C_{\chi}
$$

If we assume  $I_B=0$ , equating drain currents obtain:

$$
V_{EB2} = V_{EB1} \sqrt{\frac{W_1 L_2}{W_2 L_1}}
$$
 (7)

Thus the previous two expressions can be rewritten as :

$$
\frac{\mu C_{OX} V_{EB1}}{C_X} \left[ \frac{W_1}{L_1} \right] = \frac{\omega_0}{(\sin \theta) 2Q} \sqrt{4Q^2 - 1}
$$
(8)  

$$
\frac{\mu C_{OX} V_{EB1}}{C_X} \left[ \frac{W_1 W_2}{L_1 L_2} \right] = \frac{\omega_0}{2Q} + \frac{\omega_0}{2Q(\tan \theta)} \sqrt{4Q^2 - 1}
$$
(9)

Using the single-stage lossy integrator, design the integrator to meet a given  $\omega_0$  and Q requirement  $V_{DD}$ 

$$
\frac{\mu C_{OX}V_{EB1}}{C_X} \left[\frac{W_1}{L_1}\right] = \frac{\omega_0}{(\sin\theta)2Q} \sqrt{4Q^2-1}
$$

$$
\frac{\mu C_{OX}V_{EB1}}{C_X} \left[\frac{W_1W_2}{L_1L_2}\right] = \frac{\omega_0}{2Q} + \frac{\omega_0}{2Q(\tan\theta)} \sqrt{4Q^2-1}
$$

Taking the ratio of these two equations we obtain:

$$
\frac{W_2}{L_2} = \frac{\sin\theta + \cos\theta \sqrt{4Q^2 - 1}}{\sqrt{4Q^2 - 1}}
$$
(10)

![](_page_41_Figure_5.jpeg)

$$
I_0 = g_{m1}/C_{\chi}
$$

Observe that the pole Q is determined by the dimensions of the lossy device !

Using the single-stage lossy integrator, design the integrator to meet a given  $\omega_0$  and Q requirement  $V_{DD}$ 

$$
\frac{\mu C_{OX} V_{EB1}}{C_X} \left[ \frac{W_1}{L_1} \right] = \frac{W_0}{(\sin \theta) 2Q} \sqrt{4Q^2 - 1}
$$
\n
$$
\frac{W_2}{L_2} = \frac{\sin \theta + \cos \theta \sqrt{4Q^2 - 1}}{\sqrt{4Q^2 - 1}}
$$
\n(10)

![](_page_42_Figure_3.jpeg)

(8)

(10)

Still must obtain W<sub>1</sub>/ L<sub>1</sub>, V<sub>EB1</sub>, and C<sub>x</sub> from either of these equations

Although it appears that there might be 3 degrees of freedom left and only one constraint (one of these equations), if these integrators are connected in a loop, the operating point (Q-point) will be the same for all stages and will be that value where  $V_{out} = V_{in}$ . So, this adds a second constraint.

Setting  $V_{out}=V_{in}$ , and assuming  $V_{T1}=V_{T2}$ , we obtain from KVL

 $V_{DD} = V_{EB1} + V_{EB2} + 2V_{T}$  (11)

But  $V_{EB1}$  and  $V_{EB2}$  are also related in (7)

Using the single-stage lossy integrator, design the integrator to meet a given  $\omega_0$  and Q requirement  $V_{DD}$ 

$$
\frac{\mu C_{OX} V_{EB1}}{C_X} \left[ \frac{W_1}{L_1} \right] = \frac{W_0}{(\sin \theta) 2Q} \sqrt{4Q^2 - 1}
$$
\n
$$
\begin{array}{c}\nW_2 = \frac{\sin \theta + \cos \theta \sqrt{4Q^2 - 1}}{\sqrt{4Q^2 - 1}} \\
\frac{1}{\cos \theta} \\
\frac{1}{\
$$

Still must obtain W<sub>1</sub>/ L<sub>1</sub>, V<sub>EB1</sub>, and C<sub>x</sub> from either of these equations

$$
V_{DD} = V_{EB1} + V_{EB2} + 2V_{T}
$$
\n(11)\n
$$
V_{EB2} = V_{EB1} \sqrt{\frac{W_{1}L_{2}}{W_{2}L_{1}}}
$$
\n(12)\n
$$
V_{EB1} = \frac{V_{DD} - 2V_{T}}{1 + \sqrt{\frac{W_{2}L_{1}}{W_{1}L_{2}}}}
$$
\n(12)

 $\mathsf{I}_{\mathsf{B}}(\widehat{\mathsf{I}})$   $\Box$   $\Box$   $\mathsf{M}_{2}$ 

Substituting (10) into (12) and then into (8) we obtain

$$
\frac{\mu C_{OX}}{C_{X}} \left[ \frac{w_{1}}{L_{1}} \right] \left( \frac{V_{DD} - 2V_{T}}{1 + \sqrt{\left(\frac{W_{1}}{L_{1}}\right)^{-1} \left(\frac{\sin \theta + \cos \theta \sqrt{4Q^{2} - 1}}{\sqrt{4Q^{2} - 1}}\right)}} \right) = \frac{\omega_{0}}{(\sin \theta) 2Q} \sqrt{4Q^{2} - 1}
$$
(13)

Using the single-stage lossy integrator, design the integrator to meet a given  $\omega_0$  and Q requirement

![](_page_44_Figure_2.jpeg)

There is still one degree of freedom remaining. Can either pick  $\mathbf{W}_1\mathbf{\bm\Lambda}_1$  and solve for  $\mathbf{C}_{\mathbf{X}}$  or pick  $\mathbf{C}_{\mathbf{X}}$  and **solve for W<sup>1</sup> /L<sup>1</sup> .** 

**Explicit expression for W<sup>1</sup> /L<sup>1</sup> not available**

**Tradeoffs between C<sup>X</sup> and W<sup>1</sup> /L<sup>1</sup> will often be made**

 $\bf{S}$ ince  $\bf{V}_{\rm{OUTQ}}$ = $\bf{V}_{\rm{T}}$ + $\bf{V}_{\rm{EB1}}$ , it may be preferred to pick  $\bf{V}_{\rm{EB1}}$ , then solve (12) for  $\bf{W}_{\rm{1}}$ / $\bf{L}_{\rm{1}}$  and then solve (13) for  $C$ **x** 

Adding I<sub>B</sub> will provide one additional degree of freedom (we arbitrarily set it to 0 in this analysis) **and will relax the relationship between VOUTQ and W<sup>1</sup> /L<sup>1</sup> since (7) will be modified**

![](_page_45_Picture_0.jpeg)

# Stay Safe and Stay Healthy !

# End of Lecture 36